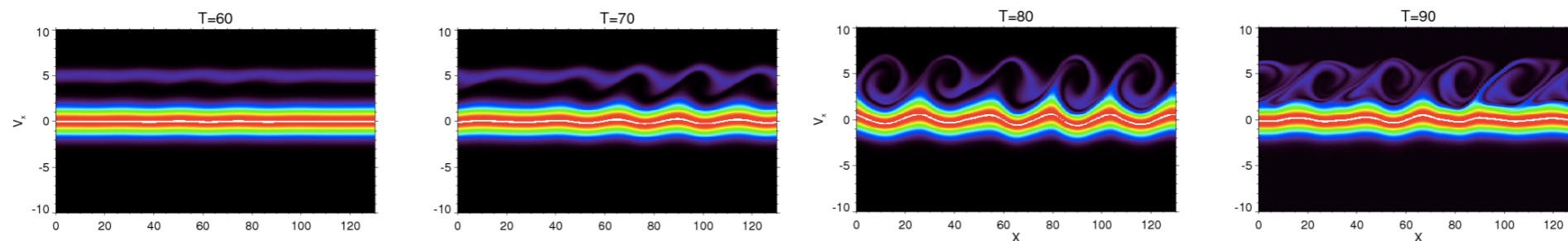


# Connecting Langmuir and low frequency waves

Pierre Henri <sup>(1)</sup> & Carine Briand <sup>(1)</sup>

(1) LPC2E, CNRS, Orléans

(2) LESIA, Observatoire de Paris



CIAS - Turbulence - May 2015

# How to transfer energy between high and low frequencies

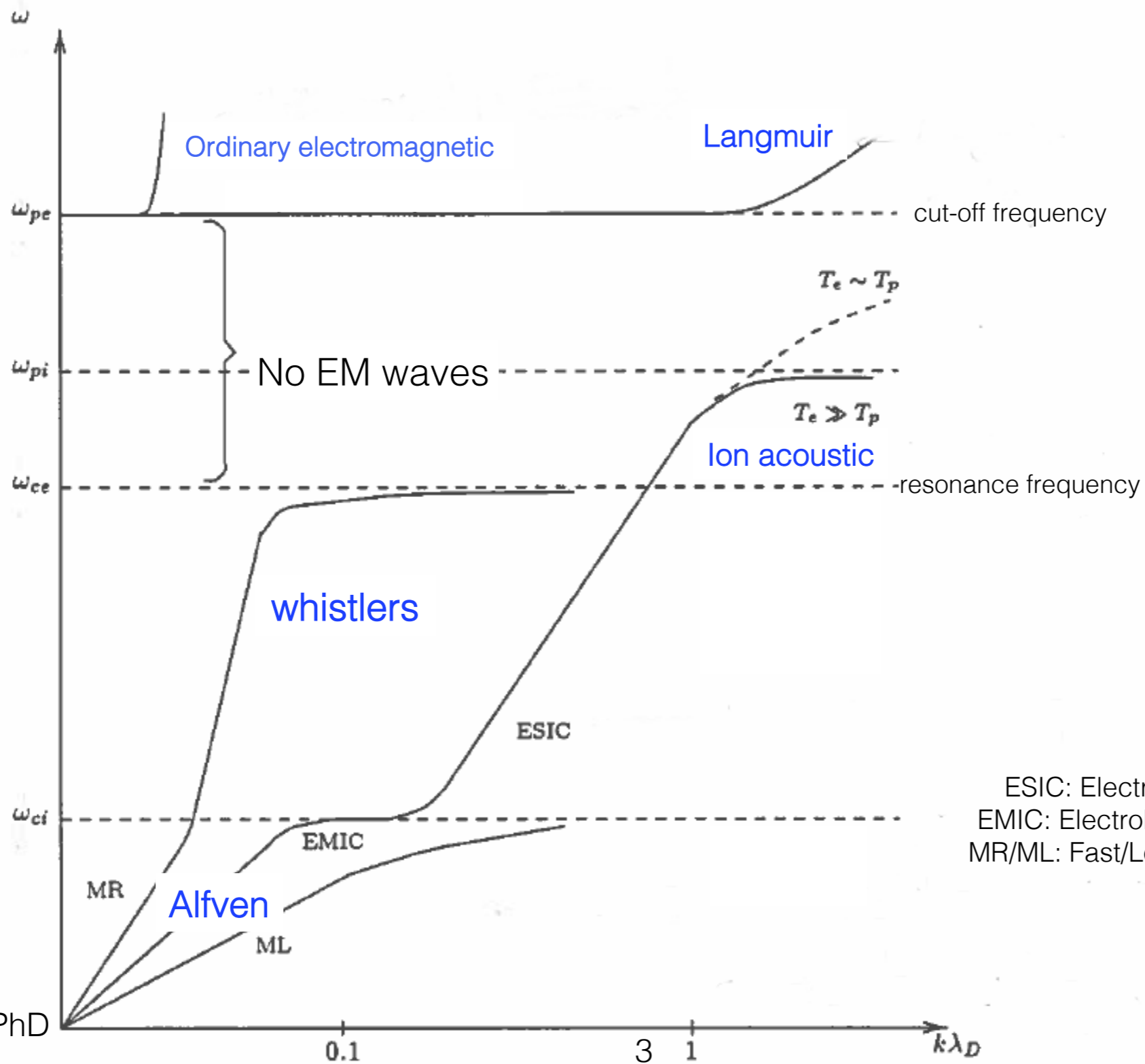
Weak turbulence

Strong Turbulence

Quasilinear evolution  
Wave-Wave interaction

Large amplitude waves / Electrons trapping

# Wave coupling



# Langmuir and ion acoustic waves

$$L \rightarrow L' + IAW$$

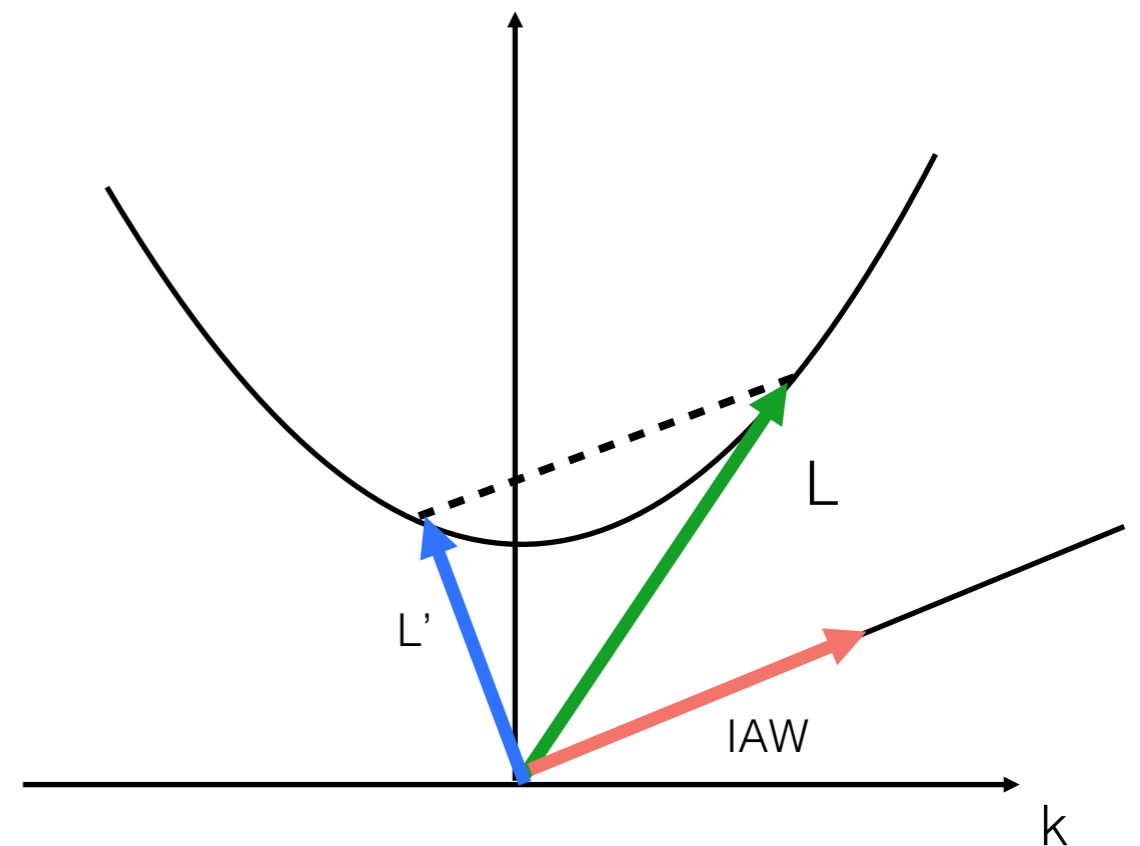
$$L + L' \rightarrow T_{2f}$$



## Resonant interaction:

- (i) conservation of momentum:  $\mathbf{k}_L = \mathbf{k}_{L'} + \mathbf{k}_s$
- (ii) conservation of energy:  $\omega_L = \omega_{L'} + \omega_s$
- (iii) phase locking:  $\varphi_L = \varphi_{L'} + \varphi_s$

ALL relations that must be checked simultaneously to verify the process



Henri, Briand, Mangeney, Bale, Califano, Goetz, Kaiser, JGR, 2009

Henri, Califano, Briand, Mangeney, JGR, 2010

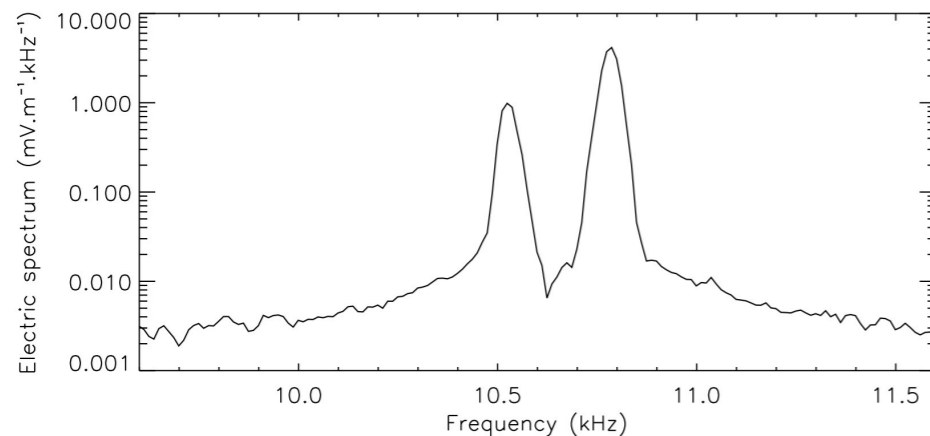
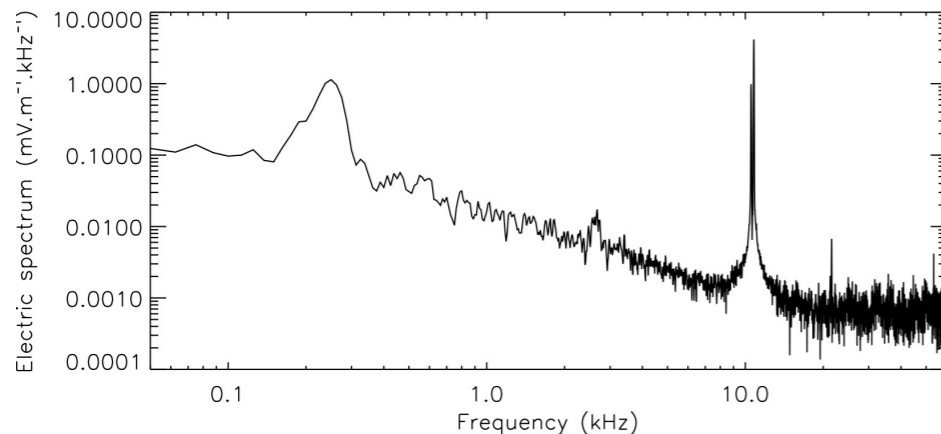
Briand, Henri, Hoang, JGR, 2014

## Before STEREO

- **spectrum:** simultaneous observations of high ( $\sim$ fp) and low ( $\sim$ 1kHz) frequency waves
- **waveform:** only access to the high frequency (too short time waveforms) - only two axes

*phase information lost*

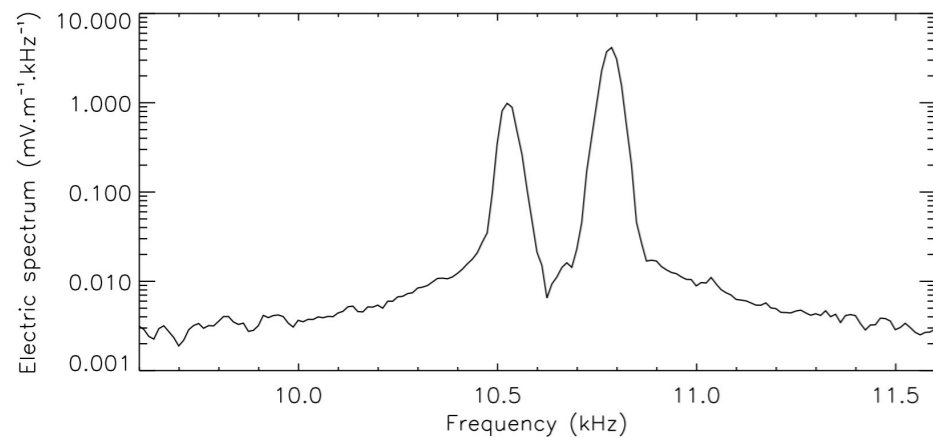
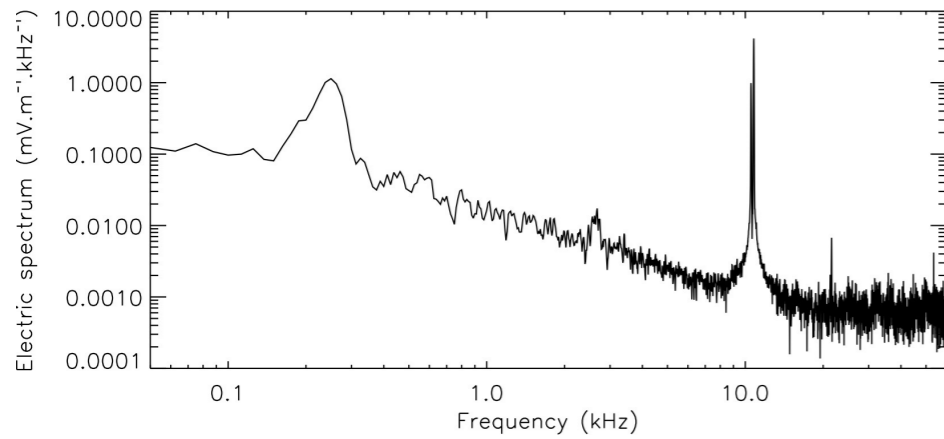
*phase information*



From STEREO waveforms we can access

- **simultaneously** to the high & low frequency signals
- 3 antennas (**3D**)

# Check of the fundamental laws

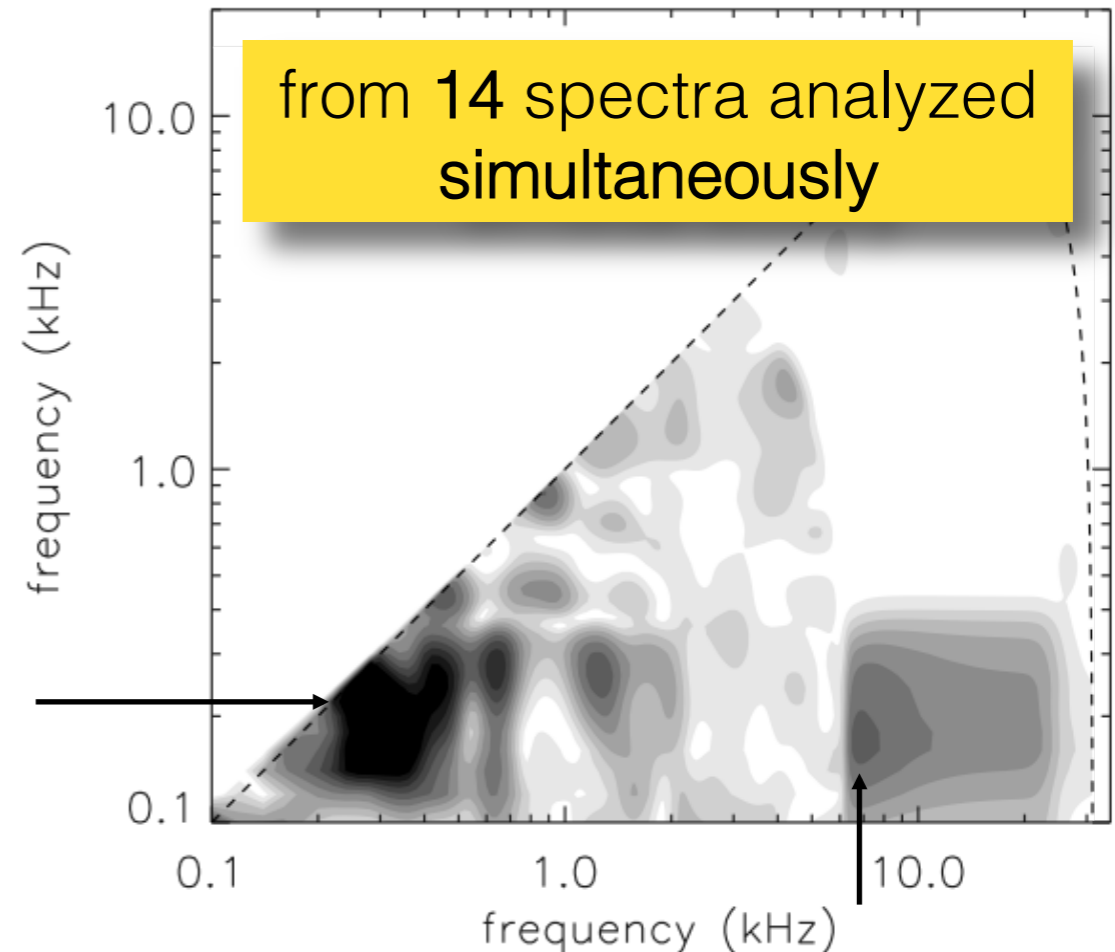


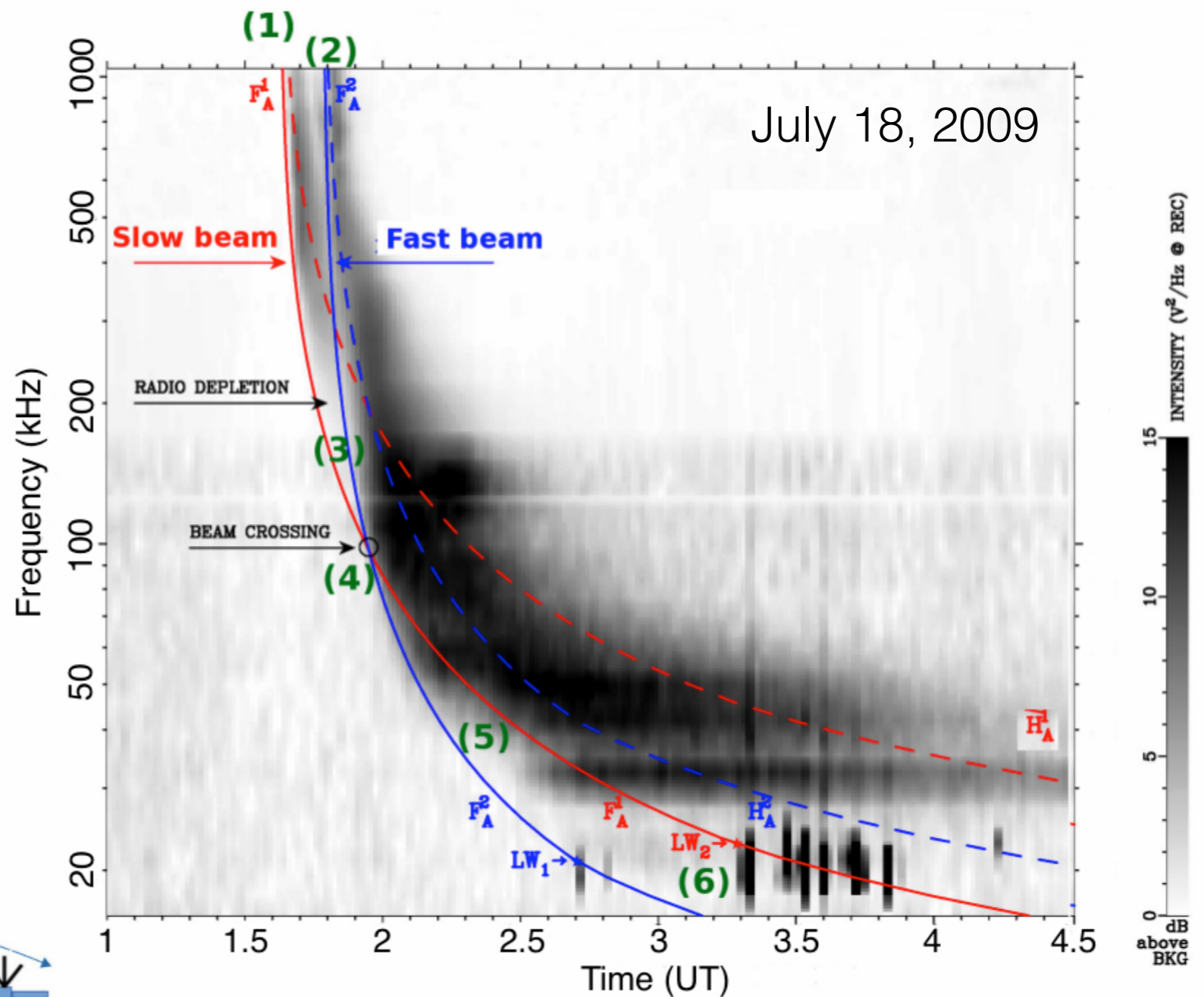
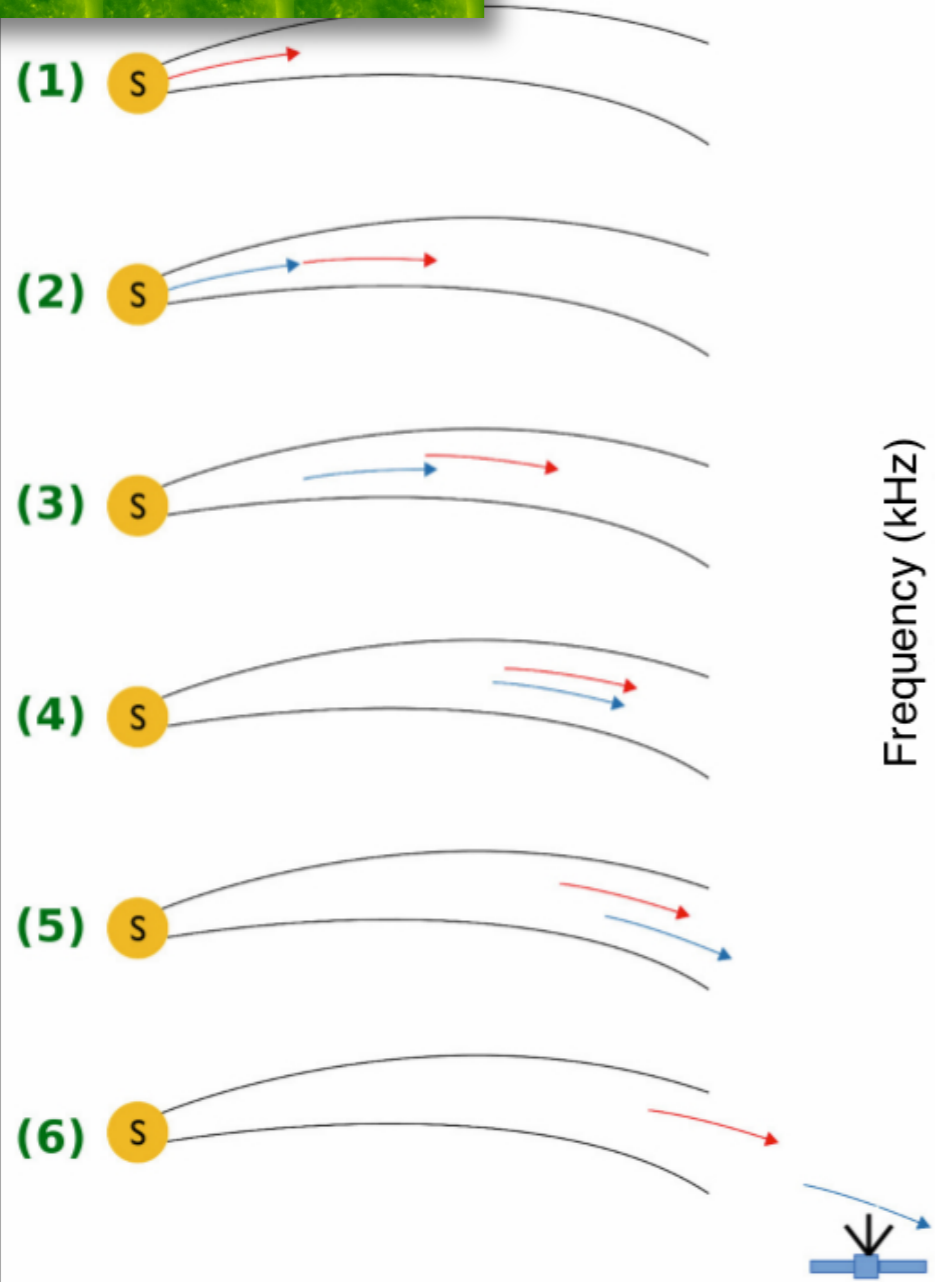
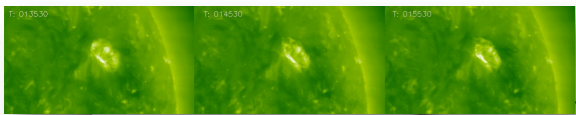
$$f_L - f_{L'} = f_{IAW}$$

taking into account Doppler-shift  
(which combines the conservation  
of energy and momentum)

## Phase locking: bicoherence

- measures the stationarity of the relative phase of the waves
- **statistical method**: to be applied to a large set of waveforms for statistical significance



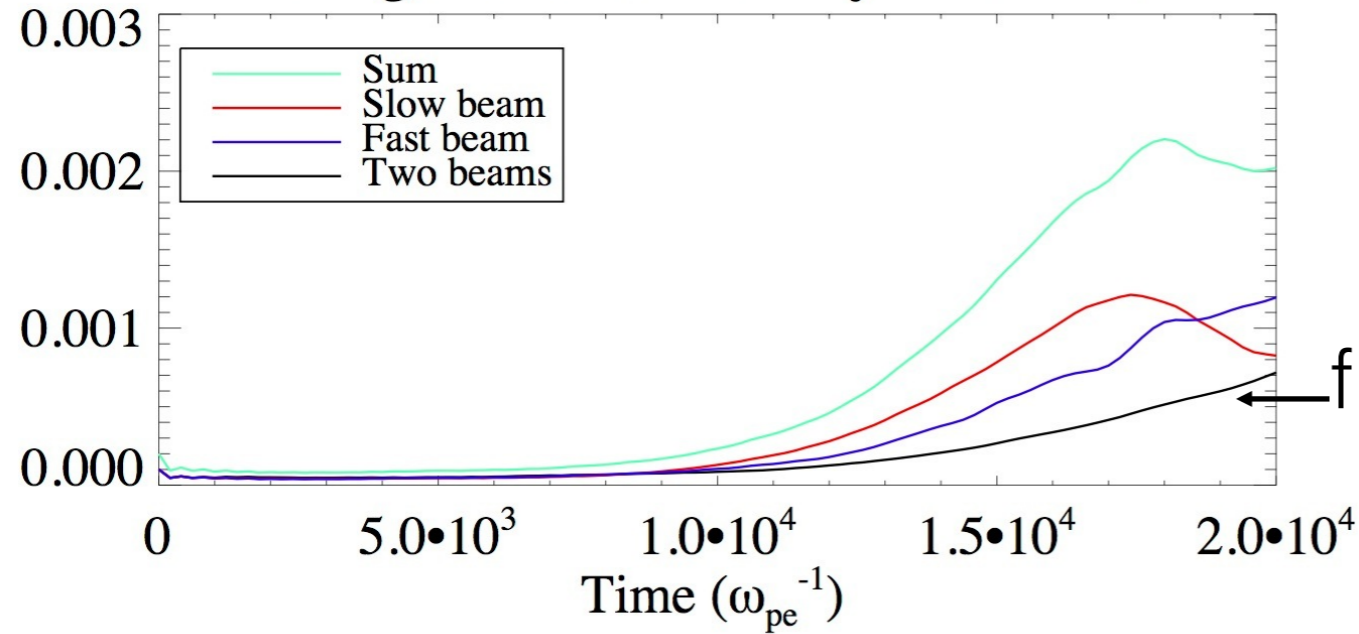


- 1° Start of a first beam (01:38) - LW detected at 03:17; slow beam (0.09c)
- 2° Start of a second beam (01:47) - LW detected at 02:42; fast beam (0.16c)

When the two beams intersect (4), reduction of the intensity of the Type III emission



## Average level of density fluctuations



With two beams, the level of density fluctuations is *always* smaller than the contribution with individual beams

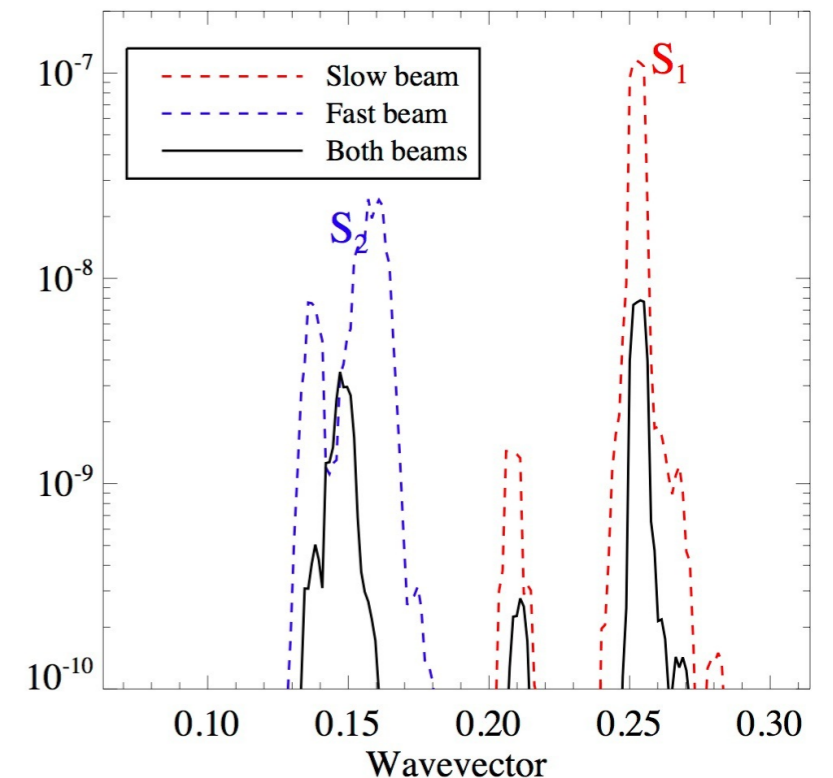
⇒ Not only the fast-beam-driven LW are affected by the presence of the slow beam.  
**slow-beam-driven LW generation is also disturbed**

## 1° Parametric process

even a small decrease of the LW energy to make the process inefficient

## 2° Coherent process

Phase beating between 3 waves. Ease if  $k_{IAW} \neq k_{L,L'}$ .  
 With 2 beams, larger spectrum of  $k$ : reduction of the phase locking



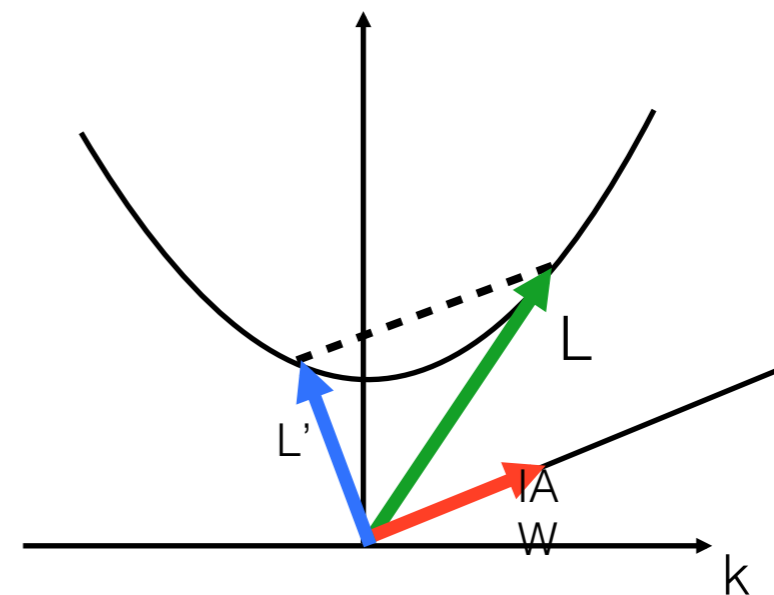
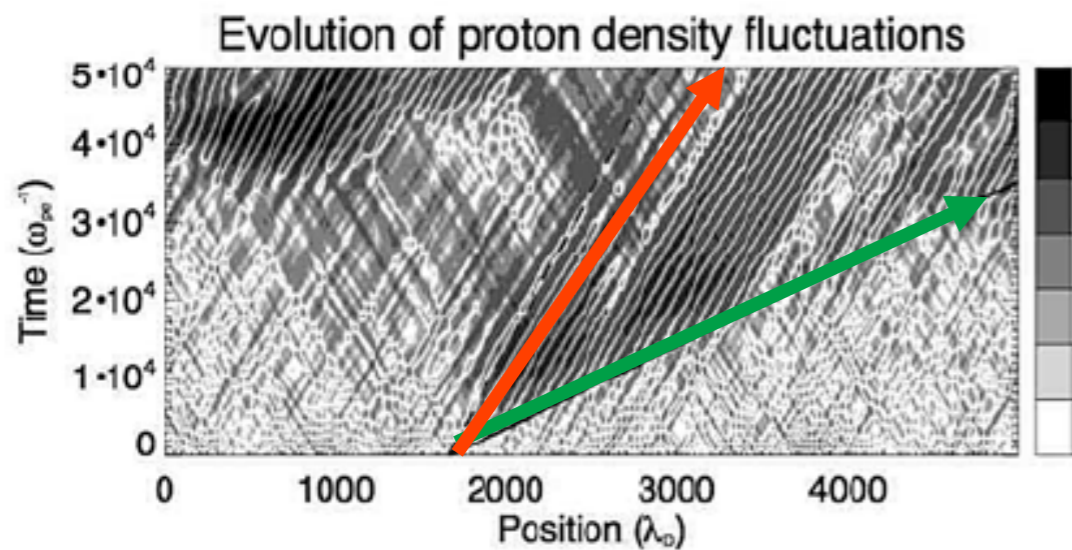
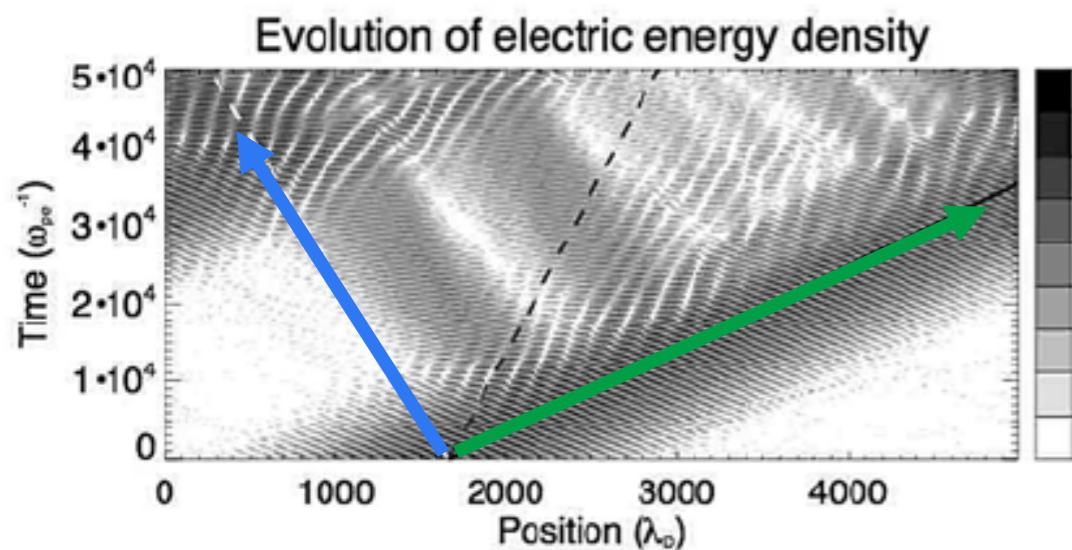


# Is the threshold for decay to proceed reached?

Vlasov-Poisson  
1D-1V electrostatic

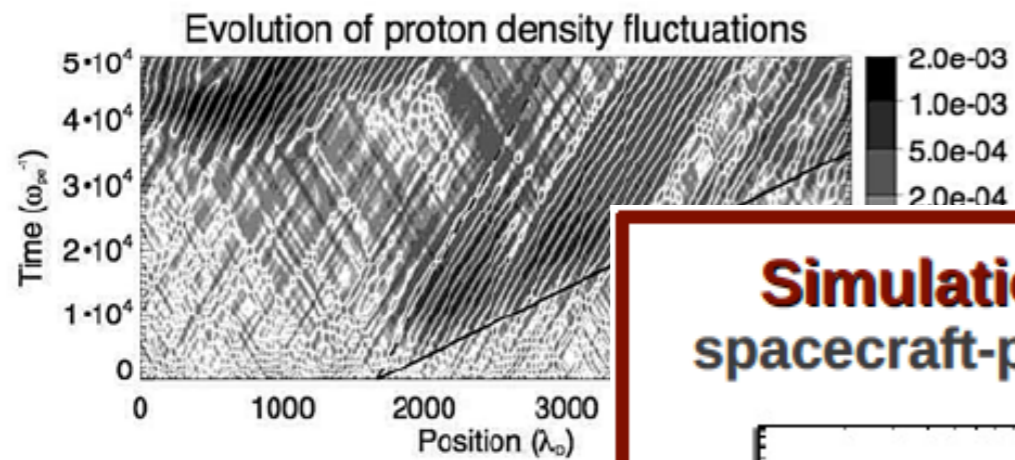
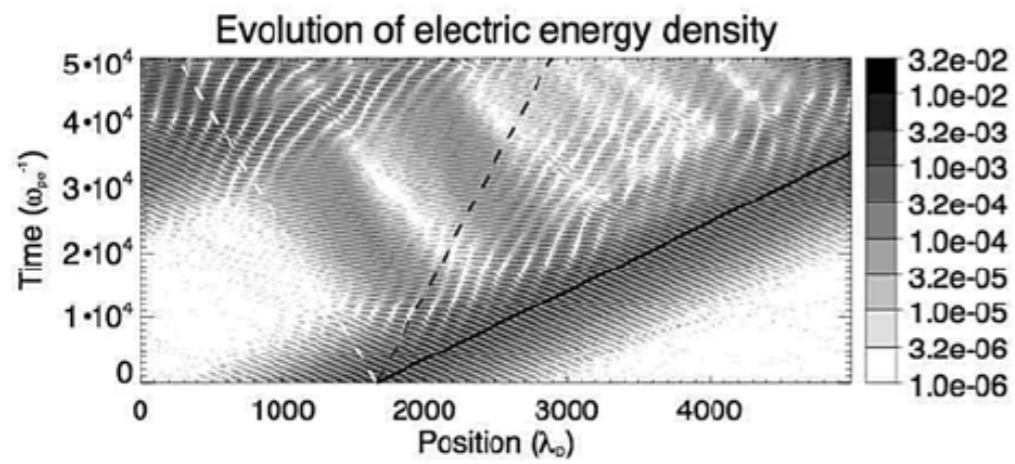
$$\frac{\partial f_\alpha}{\partial t} + v \frac{\partial f_\alpha}{\partial x} + q_\alpha E \frac{\partial f_\alpha}{\partial v} = 0$$

$$\frac{\partial E}{\partial x} = \int f_i - f_e dv$$

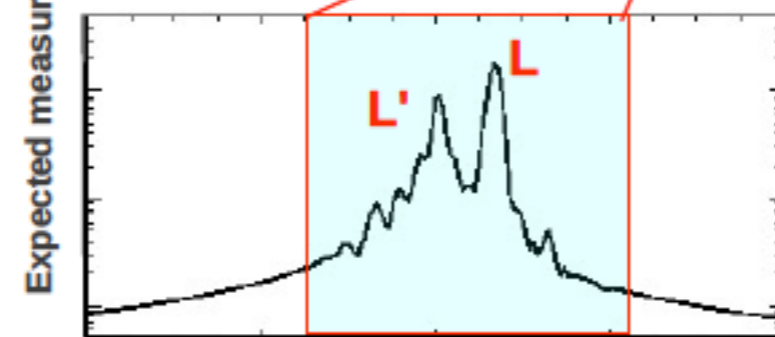
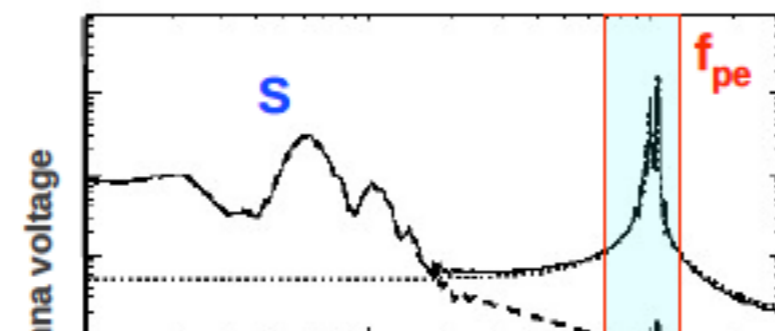


Proton density fluctuation  
(tracer of the IAW waves)

$T_e/T_i=0.1$



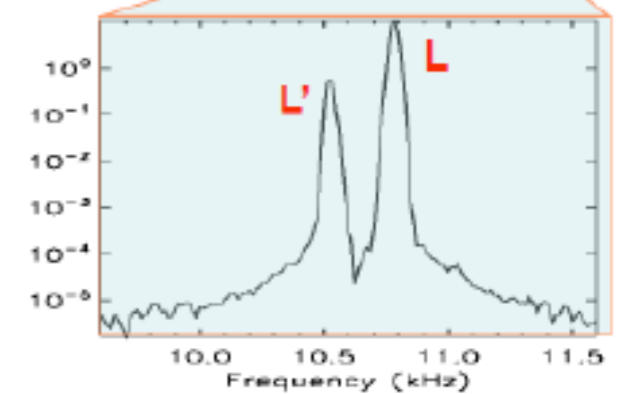
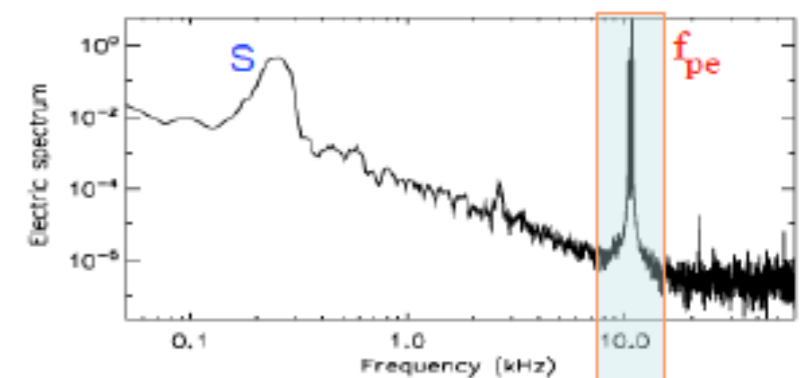
## Simulations considering spacecraft-plasma interactions



Doppler shifted frequency  
(normalised units)

## Space *in situ* observations

Fourier spectrum



[Henri et al., Solar WIND 12, 2010]

# Langmuir Electrostatic Decay: new threshold for non monochromatic waves (wave-packets of width $\Delta$ )

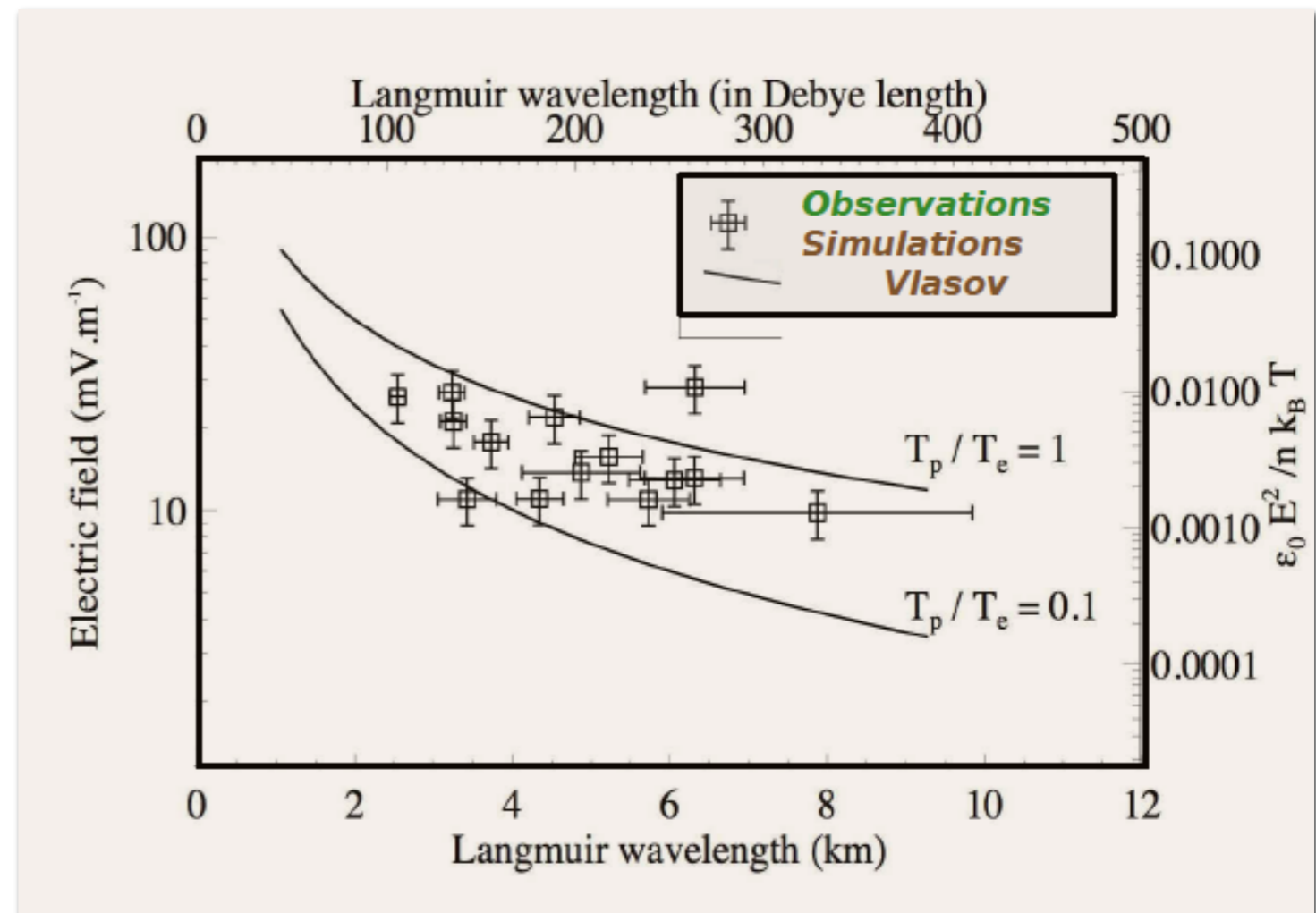
Characteristic time of wave packets collisions = instability growth rate

(function of IAW Landau damping through ion-to-electron temperature ratio)

$$E_{LED}^{thres} = \left( \frac{6k_L^{1-\beta}}{\Gamma\Delta} \right)^{1/\alpha}$$

	$\Gamma$	$\alpha$	$\beta$
$T_p/T_e = 0.1$	0.026	1.11	0.59
$T_p/T_e = 1$	0.025	1.82	0.30

Growth rate  $\gamma_{LED} = \Gamma E_L^\alpha k_L^\beta$



# Other (expected) waves for coupling

- With whistlers:
  - Radio emission: Type IV solar emission ( $L+W \rightarrow T$ ) (Kuijpers 1975; Chernov et al. 1976, 1989, 1998; Kato et al. 2014, ApJ)
  - Solar wind: Luo et al. 1999 ( $L \leftrightarrow I + W$ ; *I: left-hand circularly polarized radio wave*); Chian & Abalde 1999, Solar Phys.
  - CME: Langmuir, whistler (Moullard et al. 2012)
  - Observation of Langmuir/Upper hybrid, whistler in the auroral ionospheric region (Colpitts & Labelle 2008).
  - Separatrix of reconnection regions can lead to Langmuir, whistlers and low hybrid waves (Fujimoto, 2014);

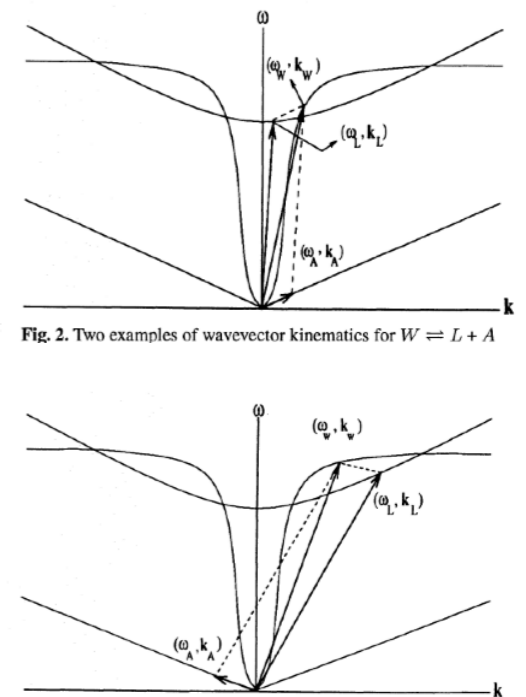
# Other (expected) waves for coupling

With kinetic Alfvén waves:  $L \leftrightarrow W + \text{KAW}$

In solar corona  
(Voitenko et al. 2003)

Planetary magnetosphere: direct coupling between Langmuir, Alfvén and whistlers (Chian, 1995, Lopes & Chian 1996)

Other coupling also possible





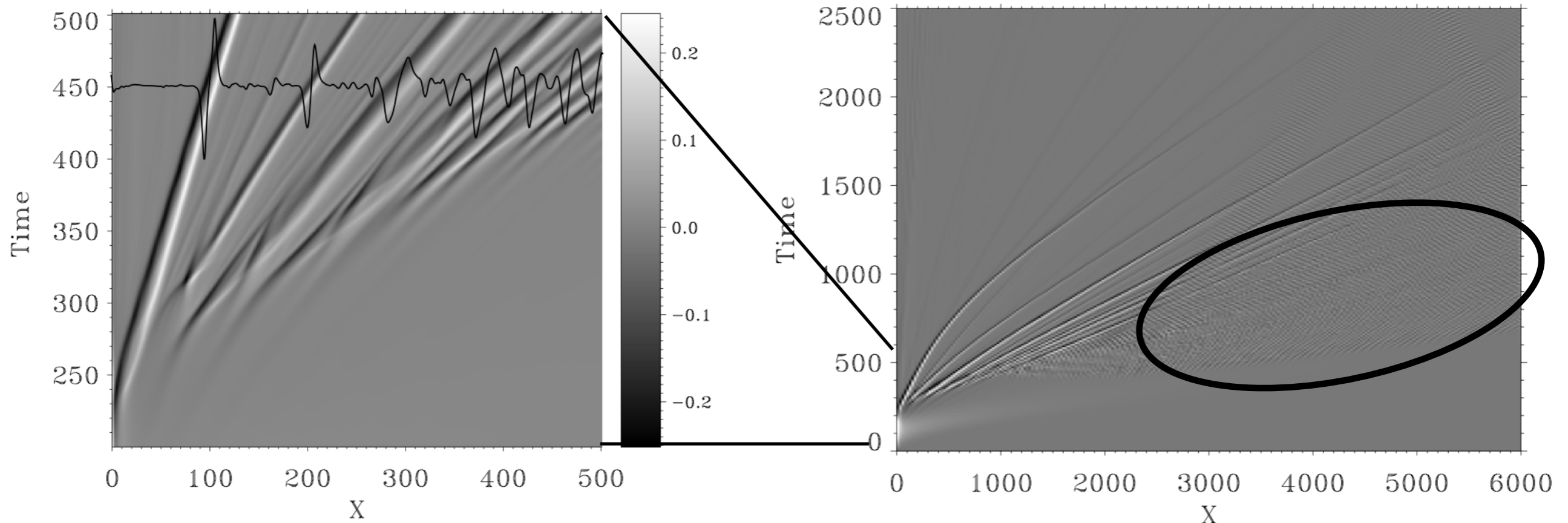
# Langmuir Turbulence: a new mechanism of generation

Temporal variation of the temperature of the electron distribution function

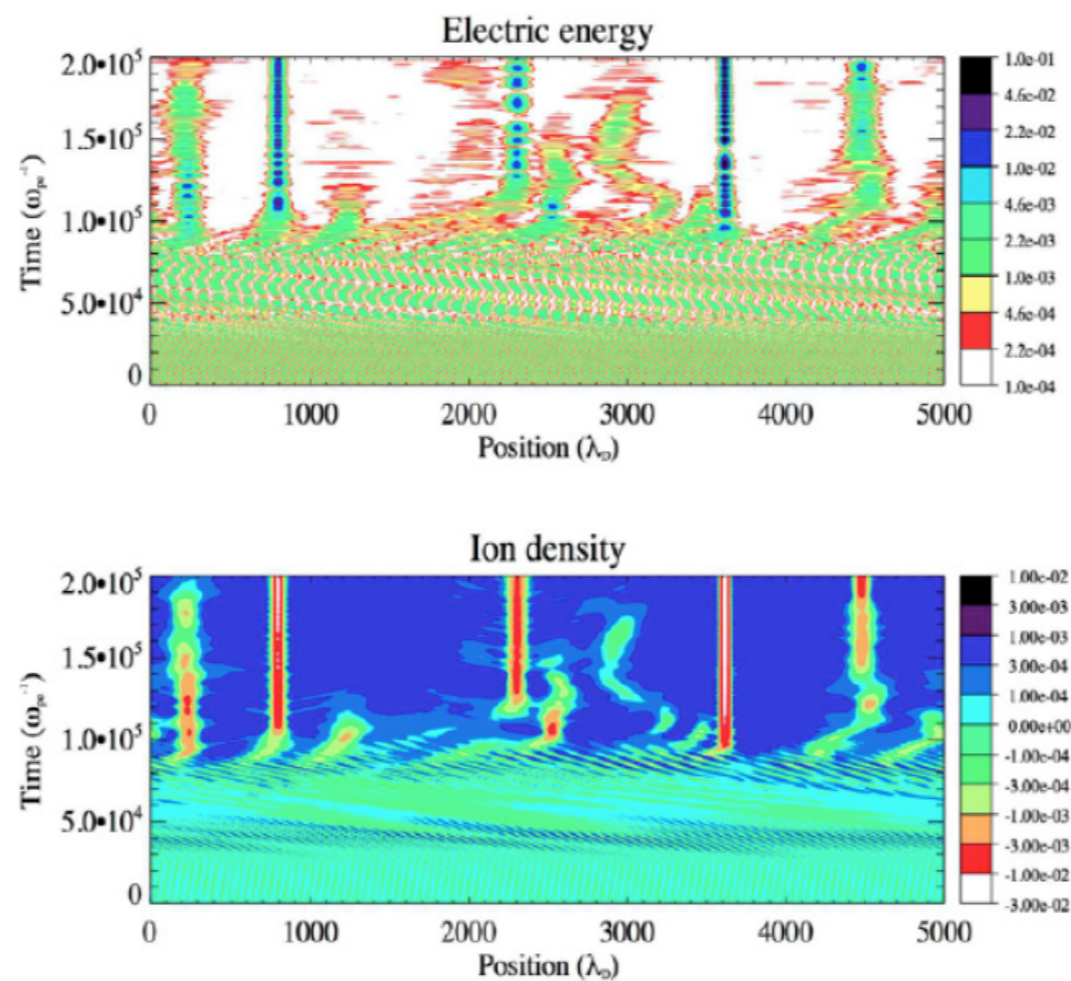
$$f_e(x=0, v_e, t) = \exp\left(-\frac{v_e^2}{2v_{\text{driver}}^2}\right),$$

$$v_{\text{driver}} = 1 + \alpha(t) \left[1 - \cos\left(\frac{2\pi t}{P}\right)\right].$$

- Development of bipolar structures
- Development of a large area of Langmuir turbulence



# C - From weak to strong Langmuir turbulence



*Henri, Califano, Briand, EPL, 2012*



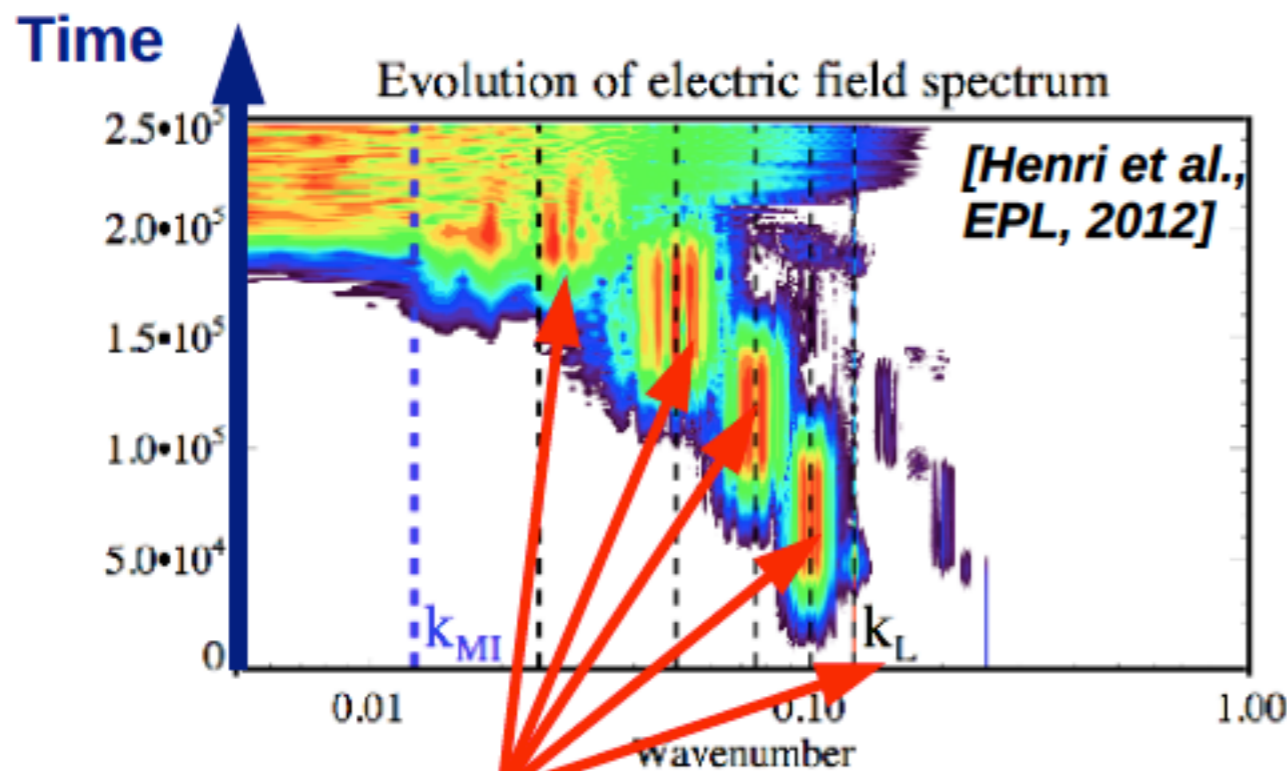
# Strong turbulence

- Collapse of soliton
- Trapping of Langmuir waves
- At the end of the collapse, transfer of the energy of the waves to the electrons.
- AGAINST:
  - In the solar wind, energy of the waves too small for strong turbulence to occur
- PRO
  - Terrestrial auroral magnetosphere/ionosphere (Eliasson 2015; Isham et al. 2012)
  - Pulsar (nanopulses) (Asseo 2006, Hanskins 2003)

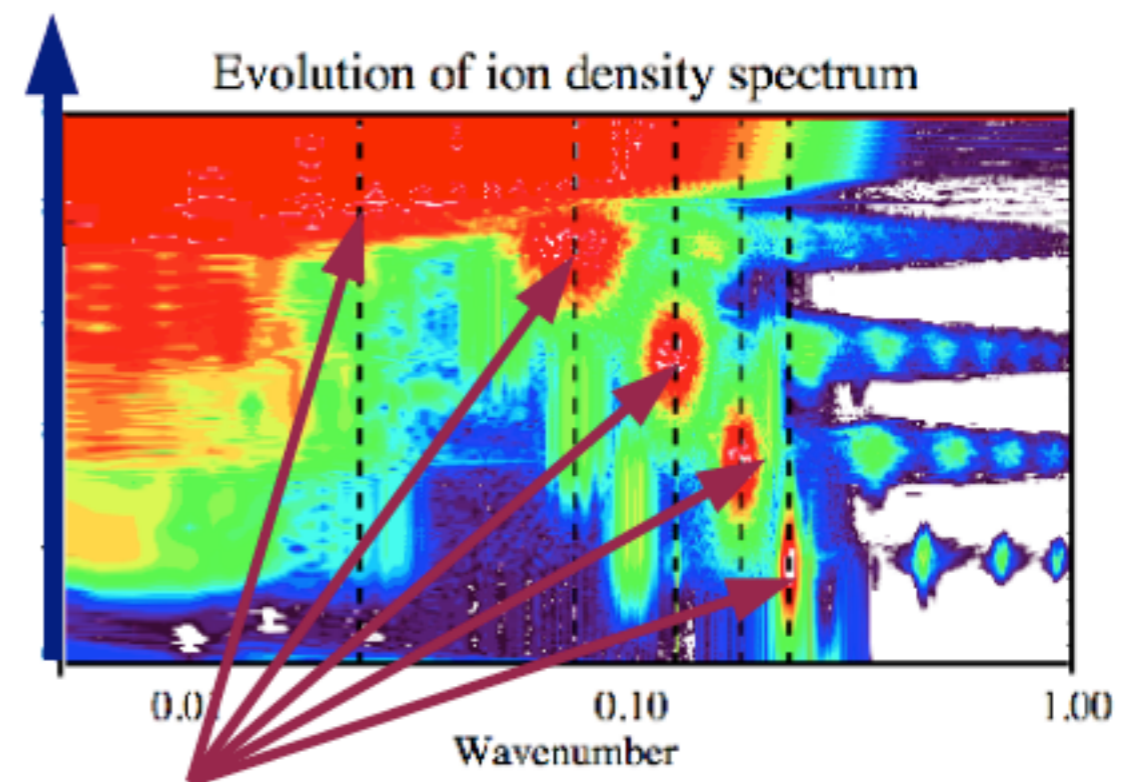
Start from a large amplitude monochromatic Langmuir waves

1st part of the evolution:

Electrostatic cascade: succession of Langmuir decay  
quasi-linear, WEAK TURBULENCE



Langmuir waves

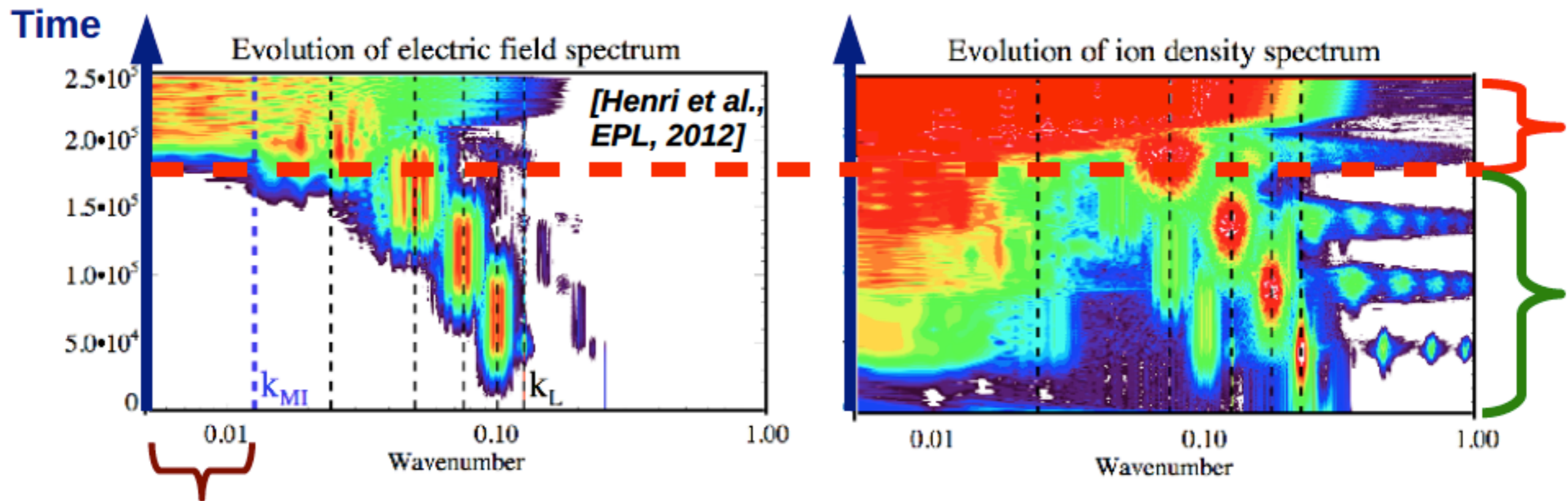


Ion Acoustic waves  
PSEUDO-MODE

@ Te=Ti

## 2nd part of the evolution:

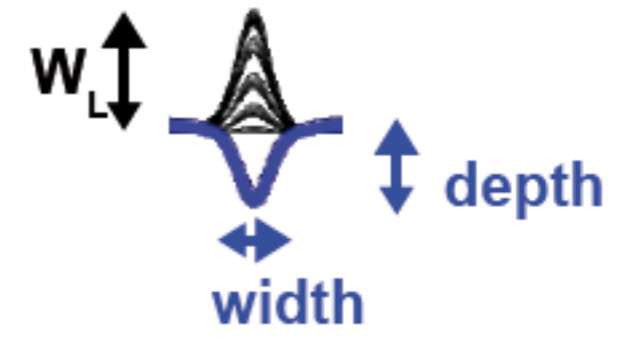
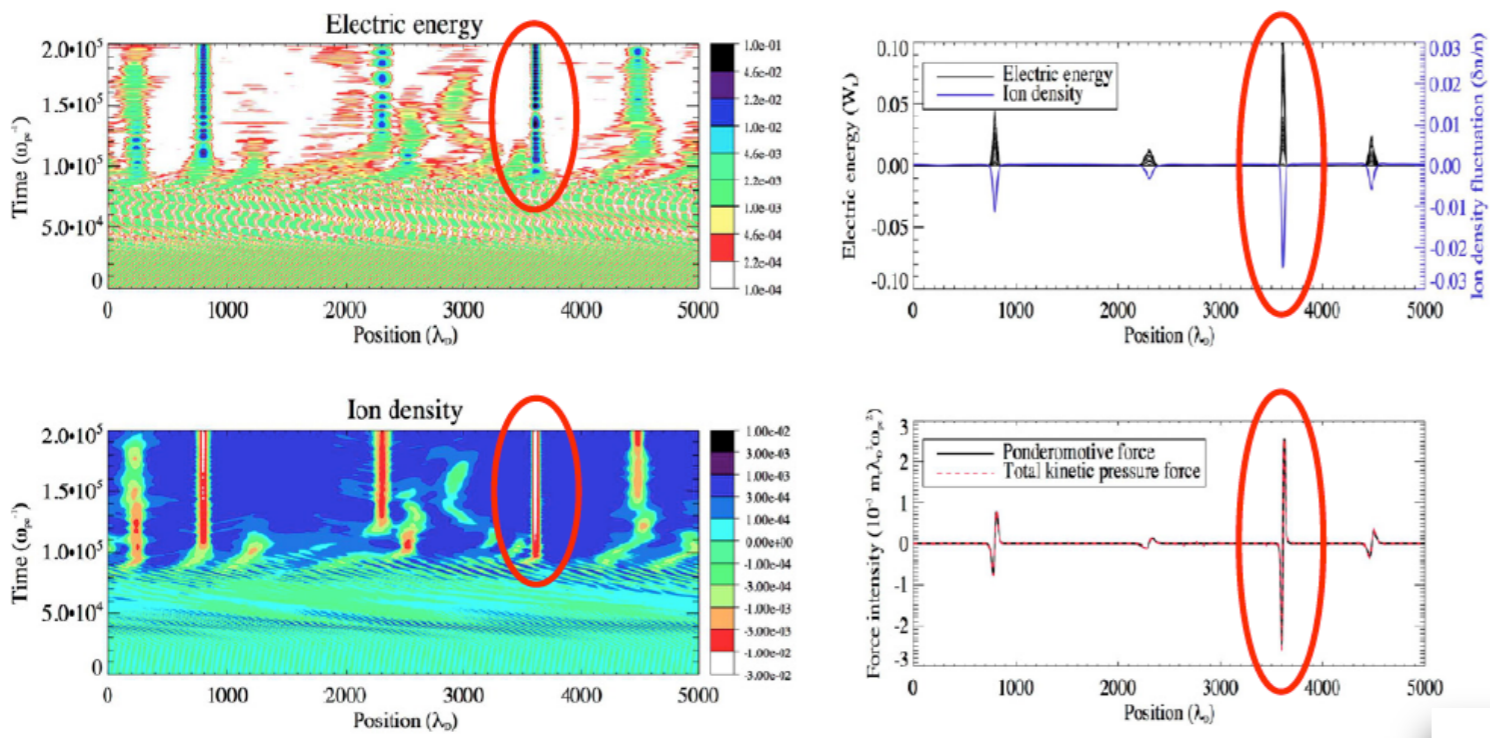
Formation of Langmuir solitons, standing coherent structures  
strongly nonlinear, STRONG turbulence regime



Langmuir modulational instability

strong turbulence  
weak turbulence

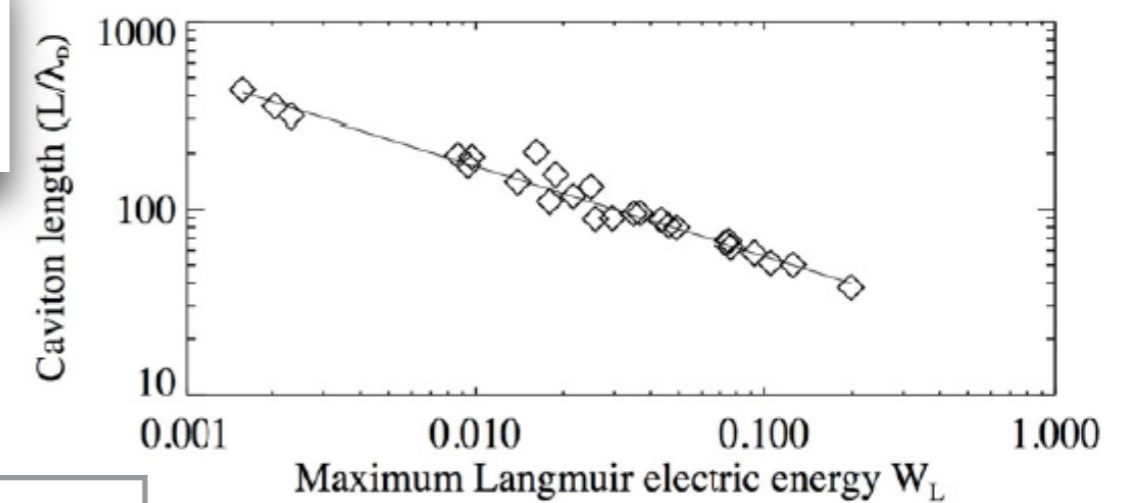
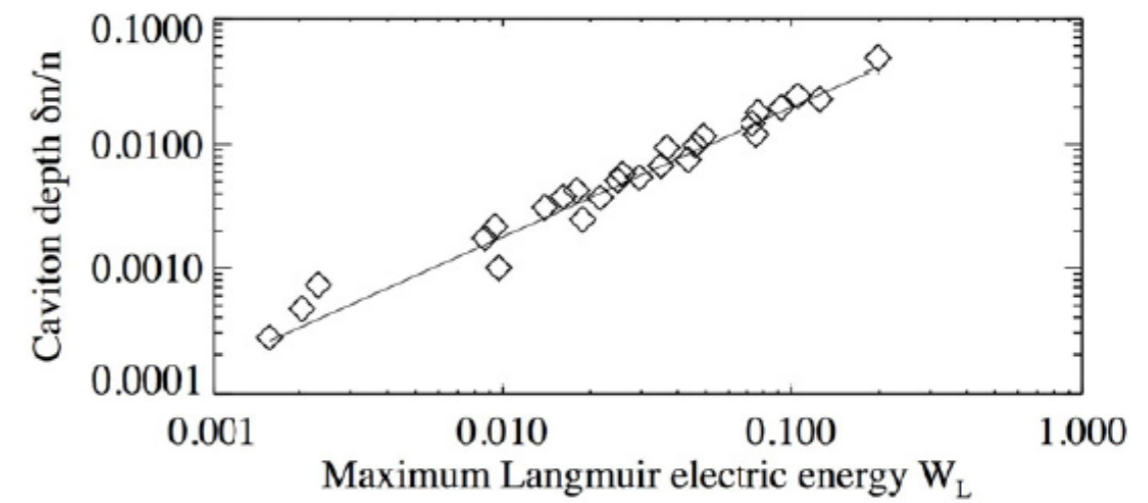




Example of (standing) Langmuir Caviton

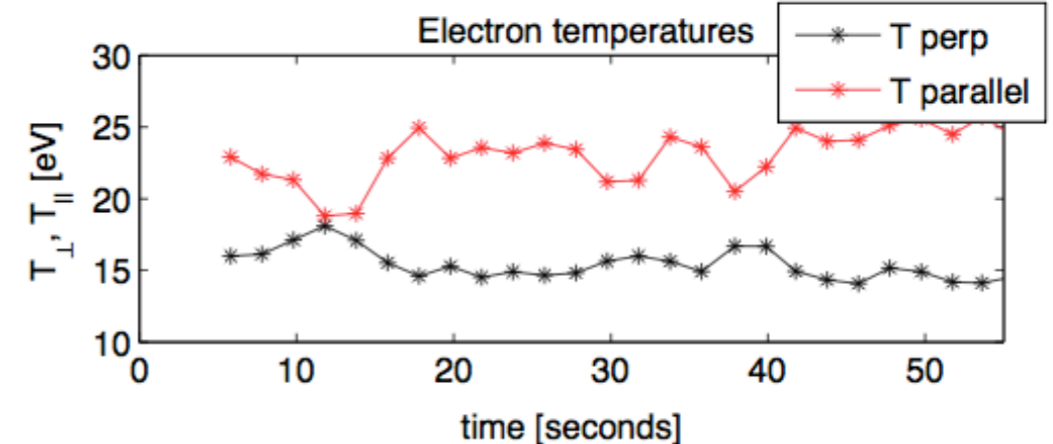
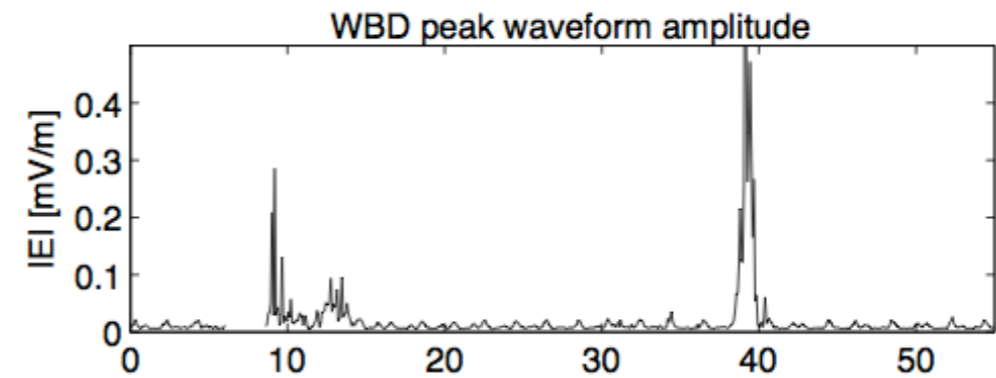
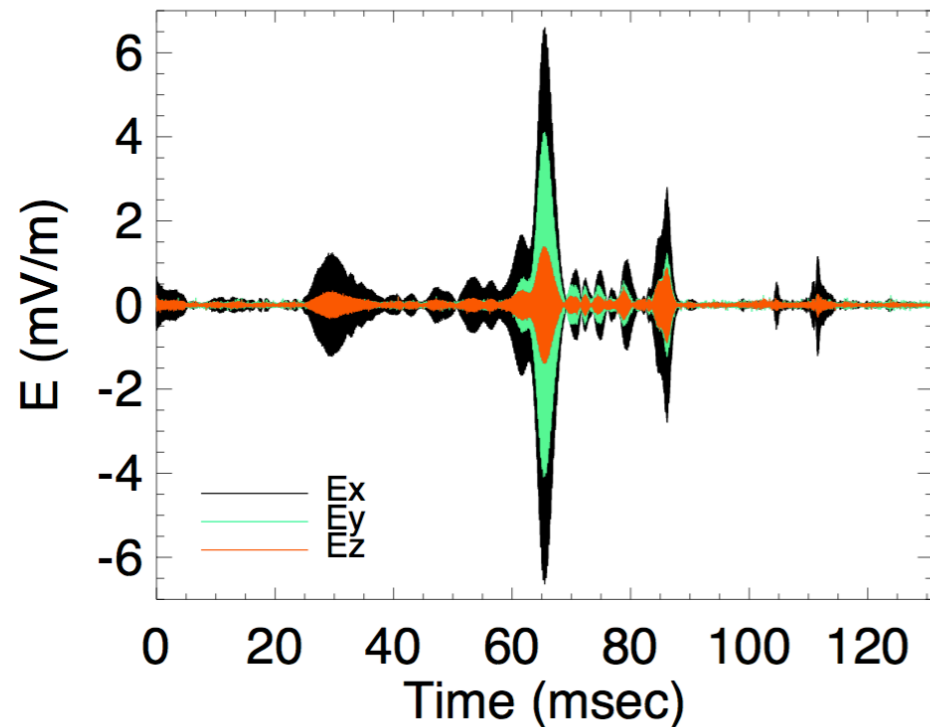
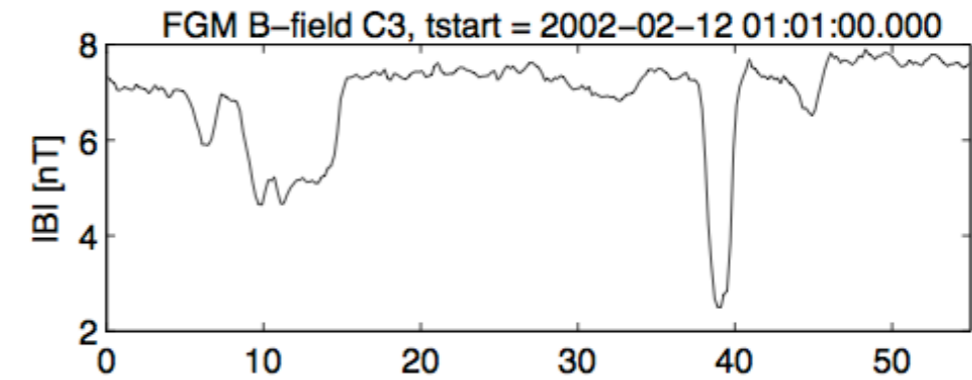
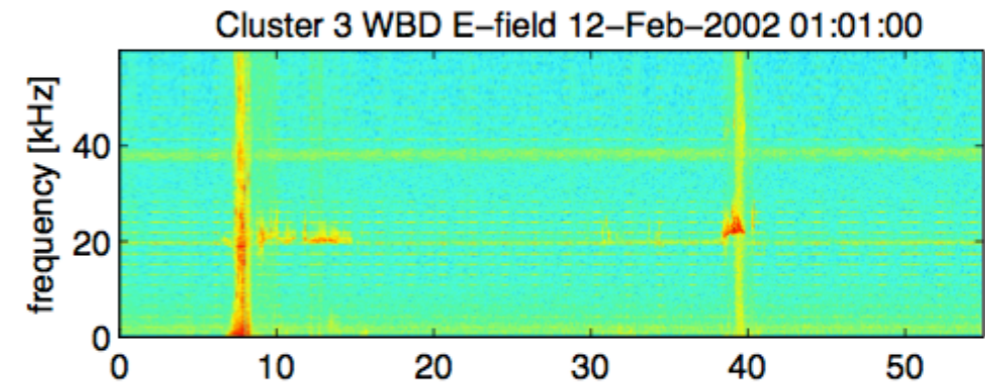
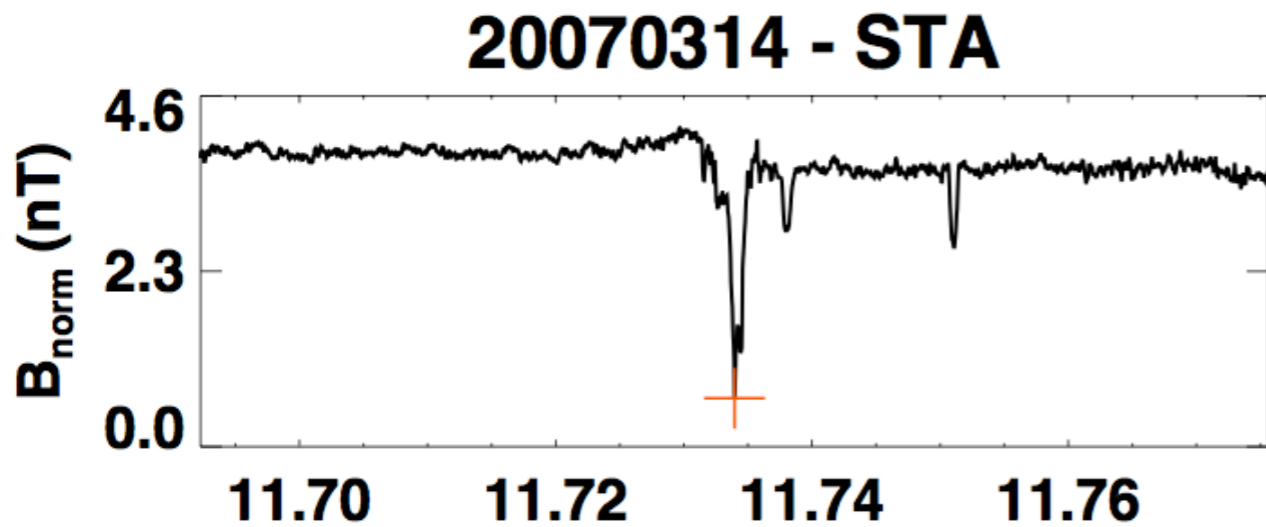
Same scaling laws for *standing* and *propagating* cavitons (Langmuir solitons)

### Scaling laws for (standing) Langmuir Cavitons



**Cavitons can saturate at low energy (until  $\epsilon_0 E^2 / nk_B T \sim 10^{-3}$ ) unlike the wide-spread belief that such structures should be found at higher energy ratios ( $\epsilon_0 E^2 / nk_B T \sim 1$ ).**

# Wave trapping in magnetic holes



# Summary & futur works

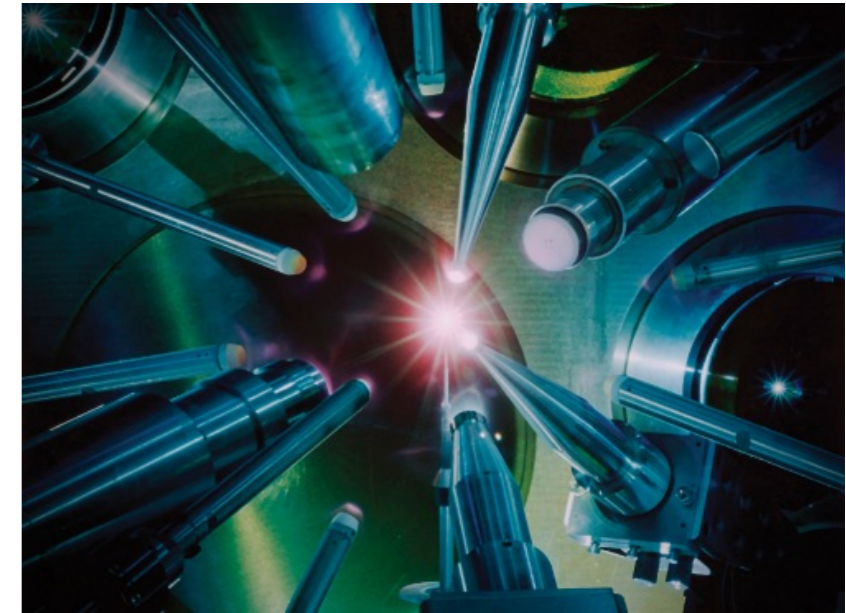
# Observing requisite

- **Waveforms are mandatory** to check the conditions for wave coupling ( $k$ ,  $w$ , phase locking and coherence)
- Futur instruments must be able to provide electric waveforms to detect from a few Hz to hundred(s) kHz to study the different coupling and transfer of energy from high to low frequencies
- Particle Distribution Function: high resolution in **energy** (some instabilities very sensitive to the thermal velocity of the beam), and high time resolution



# A universal question

Collaboration: inertial confinement fusion community (LULI)  
PIC simulations + data from laser experiments



Laser is dispersed through *Ion Acoustic Parametric Decay*

Production of hot electrons

Ion waves: source of anomalous resistivity

Saturation of the Raman instability

	$\rho_e \text{ (cm}^{-3}\text{)} / f_p$	$T_p/T_e$	$T_e \text{ (keV)}$	$k_{\perp}\lambda_d$
<b>ICF(1micron)</b>	$10^{12-19} / 900\text{MHz}$	0.1-1	0.1-1	0.01-0.5
<b>Space (1AU)</b>	$1-10 / 10 - 30\text{kHz}$	0.1 - 1	$10^2$	< 0.05